

# Periodic optimal control, dissipativity and MPC

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**Abstract**—Recent research has established the importance of (strict) dissipativity for proving stability of economic MPC in the case of an optimal steady state. In many cases, though, steady state operation is not economically optimal and periodic operation of the system yields a better performance. In this paper, we propose ways of extending the notion of (strict) dissipativity for periodic systems. We prove that optimal  $P$ -periodic operation and MPC stability directly follow, similarly to the steady state case, which can be seen as a special case of the proposed framework. Finally, we illustrate the theoretical results with several simple examples.

**Index Terms**—Periodic Economic Model Predictive Control, Strict dissipativity

## I. INTRODUCTION

Economic MPC is a variant of model predictive control (MPC) in which the objective consists in directly optimizing a given performance index as opposed to tracking a given reference. The main advantage of economic MPC over tracking MPC becomes apparent in transients, when the system is steered to steady state while minimizing the given performance index.

Unfortunately, proving stability of economic MPC schemes is hard, as the stage cost  $\ell(x, u)$  does in general not have a pointwise minimum on the trajectory the system converges to. The idea of rotating the cost using the Lagrange multipliers  $\lambda$  has been proposed in [4] in order to prove stability. The proof relies on an equivalent auxiliary MPC scheme with a rotated stage cost that has a stationary point at the optimal steady state. The rotated stage cost is obtained by adding the term  $\lambda^T x - \lambda^T f(x, u)$  to the stage cost. In [1] this idea has been extended to a nonlinear rotation, given by a function  $\lambda(x)$ . This generalization is equivalent to the systems theoretic notion of strict dissipativity [11], [12] with  $\lambda$  as a storage function and allows one to both rotate and lower bound the stage cost of the auxiliary MPC scheme. For a given system and stage cost, if there exists a storage function  $\lambda(x)$  that satisfies a strict dissipativity property, then stability of the MPC scheme is guaranteed.

A first extension of this framework to periodic systems has been proposed in [13] in the context of time varying systems, where the Lagrange multipliers  $\lambda_k$  of a periodic optimal trajectory have been used to rotate the cost with a linear (time varying) term. In contrast to this reference, in the present paper we consider optimal periodic trajectories for optimal control problems with time invariant dynamics and stage costs. To this end, we propose and discuss two different ways of extending the notion of dissipativity to the periodic

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case in order to rotate the stage cost of the auxiliary MPC scheme for proving optimality properties of periodic orbits and stability of periodic economic MPC schemes. For the latter, we focus on economic MPC formulations with appropriate terminal constraints and costs.

The paper is structured as follows. Section II introduces the basic notation and summarises previous results obtained for the steady state case. The newly proposed concept of  $P$ -periodic dissipativity is introduced in Section III, where previous results on optimal operation at steady state are extended to the periodic case. The stability proof for periodic economic MPC is given in Section IV. Some simple examples are presented in Section V in order to illustrate the theory. Conclusions and a discussion on future research directions are given in Section VI.

## II. SETTING

We consider discrete time nonlinear systems governed by the dynamics

$$x_{k+1} = f(x_k, u_k) \quad (1)$$

with  $f : X \times U \rightarrow X$ , with  $X \subset \mathbb{R}^n$  and  $U \subset \mathbb{R}^m$ . Solutions for initial value  $x_0$  and control sequence  $u$  are denoted by  $x_k^u(x_0)$ .

We assume that  $f$  is continuous in  $(x, u)$  and the system is subject to state and input constraints  $(x_k, u_k) \in \mathbb{Z} \subset \mathbb{X} \times \mathbb{U}$  for all  $k \geq 0$ . In the MPC framework, the system is equipped with a stage cost  $\ell : \mathbb{X} \times \mathbb{U} \rightarrow \mathbb{R}$  which is assumed to be continuous.

For given state and control constraint set  $\mathbb{Z}$ , each initial value  $x_0 \in \mathbb{X}$  and any  $N \geq 1$  we denote the admissible control sequences by  $\mathbb{U}^N(x_0) := \{u(\cdot) \in \mathbb{U}^N \mid (x_k^u(x_0), u_k) \in \mathbb{Z} \forall k = 0, \dots, N-1\}$ . Analogously we define  $\mathbb{U}^\infty(x_0)$ . For simplicity of exposition we assume  $\mathbb{Z}$  to be compact. We consider the finite horizon functional

$$J_N(x, u) := \sum_{k=0}^{N-1} \ell(x_k^u(x), u_k)$$

and the infinite horizon averaged functional

$$\bar{J}_\infty(x, u) := \limsup_{K \rightarrow \infty} \frac{1}{K} J_K(x, u).$$

which are well defined for all  $u \in \mathbb{U}^N(x)$  or  $u \in \mathbb{U}^\infty(x)$ , respectively, provided  $\ell$  is bounded along the trajectory.

Given an initial value  $x_0^{\text{MPC}} \in \mathbb{X}$ , the basic model predictive control (MPC) scheme with nominal system dynamics works as follows:

- (i) set  $n := 0$
- (ii) minimise  $J_N(x_n^{\text{MPC}}, u)$  over all control sequences  $u \in \mathbb{U}^N(x_n^{\text{MPC}})$  and denote the optimal sequence by  $u^*$
- (iii) set  $u_n^{\text{MPC}} := u_0^*$ ,  $x_{n+1}^{\text{MPC}} := f(x_n^{\text{MPC}}, u_n^{\text{MPC}})$ ,  $n := n + 1$  and go to (ii)

Since the stage cost  $\ell$  is not of tracking type (i.e., does not necessarily penalise the distance to a pre-specified equilibrium) this MPC scheme is often termed *economic MPC* [1], [2]. In this setting, the classical notion of (strict) dissipativity [11], [12] has recently gained renewed interest.

**Definition 2.1 (Strict Dissipativity [1]):** System (1) is dissipative with respect to a steady state  $(x^s, u^s) \in \mathbb{Z}$  of (1) for supply rate  $\ell(x, u) - \ell(x^s, u^s)$  if there exists a storage function  $\lambda : \mathbb{X} \rightarrow \mathbb{R}$  such that the inequality

$$L(x, u) := \ell(x, u) - \ell(x^s, u^s) + \lambda(x) - \lambda(f(x, u)) \geq 0$$

holds for all  $(x, u) \in \mathbb{Z}$ . If, in addition, there exists a function  $\rho \in \mathcal{K}_\infty$  such that the inequality

$$L(x, u) \geq \rho(\|x - x^s\|)$$

holds, then the system (1) is strictly dissipative on  $\mathbb{Z}$ .  $\square$

If a system equipped with a stage cost  $\ell$  is (strictly) dissipative, then this has several consequences:

- The system is optimally operated at (uniformly suboptimally operated off) steady state [2], [8].
- For economic MPC with terminal constraint, the averaged performance  $\bar{J}_\infty(x^{\text{MPC}}, u^{\text{MPC}})$  equals  $\ell(x^s, u^s)$  and the steady state  $x^s$  is asymptotically stable for the closed loop solutions. This was shown for periodic endpoint constraints in [4] for linear storage functions and in [2] for general storage functions as well as for regional constraints and terminal costs in [1].
- For economic MPC without terminal constraint, the averaged performance  $\bar{J}_\infty(x^{\text{MPC}}, u^{\text{MPC}})$  equals  $\ell(x^s, u^s) + \varepsilon(N)$  and the optimal equilibrium is practically asymptotically stable, cf. [7], [5]. Moreover, approximate transient optimality was shown in these references and — under an exponential turnpike property which in turn is implied by dissipativity and suitable controllability properties [3] — the error terms converge to 0 exponentially fast as  $N \rightarrow \infty$ .

For general discrete time optimal control problems, it is well known that the optimal value is not necessarily attained at an equilibrium. Particularly, it may happen that periodic orbits exhibit smaller average values than any feasible equilibrium, see, e.g., [2, Section VII] or our examples below. In this case, the existing theory based on dissipativity of an equilibrium is not applicable and does thus not ensure asymptotic stability of the optimal periodic orbit. For this reason, in the next section we discuss dissipativity notions which are adapted to characterising periodic orbits.

### III. PERIODIC DISSIPATIVITY

In this section, we will introduce concepts of  $P$ -periodic (strict) dissipativity and analyse how they relate to optimal  $P$ -periodic operation. Periodic EMPC stability will then be addressed in the following section. Let us first give definitions of periodic orbits and periodic trajectories.

*Definition 3.1 (Periodic Orbit):* An ordered  $P$ -tuple of points  $\Pi = (\bar{x}_0^p, \dots, \bar{x}_{P-1}^p)$ ,  $P \geq 1$ , is called a *feasible  $P$ -periodic orbit* with control sequence  $(\bar{u}_0^p, \dots, \bar{u}_{P-1}^p)$  if  $(\bar{x}_k^p, \bar{u}_k^p) \in \mathbb{Z}$ ,  $k = 0, \dots, P-1$ ,  $\bar{x}_0^p = \bar{x}_{P-1}^p$  and

$$\bar{x}_{k+1}^p = f(\bar{x}_k^p, \bar{u}_k^p) \quad \text{for } k = 0, \dots, P-1.$$

The number  $P$  is called the *period* of the orbit  $\Pi$  and if there is no  $Q \geq 1$  with  $Q < P$  such that  $(\bar{x}_k^p, \bar{u}_k^p) = (\bar{x}_{k+Q}^p, \bar{u}_{k+Q}^p)$  for all  $k = 0, \dots, P-Q$ , then  $P$  is called the *minimal period* of  $\Pi$ . Given the corresponding control sequence  $(\bar{u}_0^p, \dots, \bar{u}_{P-1}^p)$  we define the tuple of state-control pairs  $\Pi_U := ((\bar{x}_0^p, \bar{u}_0^p), \dots, (\bar{x}_{P-1}^p, \bar{u}_{P-1}^p))$ .  $\square$

Note that in our terminology an equilibrium is a periodic orbit with period  $P = 1$ . Moreover, for  $P > 1$ , the periodic orbit is not unique, as phase shifts produce an orbit which is defined by the same states and controls, but in a shifted order. For this reason, we define in the following the periodic trajectory as a periodic orbit with a fixed phase, extended infinitely long into the future.

*Definition 3.2 (Periodic Trajectory):* (i) A sequence  $X^P = (x_0, x_1, x_2, \dots)$ , is called a *feasible  $P$ -periodic trajectory* with control sequence  $U^P = (u_0, u_1, u_2, \dots)$  if  $(x_k, u_k) \in \mathbb{Z}$ ,  $x_k = x_{k+P}$ ,  $u_k = u_{k+P}$  for all  $k = 0, 1, \dots$ , and

$$x_{k+1} = f(x_k, u_k) \quad \text{for } k = 0, 1, \dots$$

(ii) Given a  $P$ -periodic orbit  $\Pi = (\bar{x}_0^p, \dots, \bar{x}_{P-1}^p)$  and a *phase*  $\phi \in \{0, \dots, P-1\}$ , we define the infinite sequence

$$X_\phi^P(\Pi) := (\bar{x}_\phi^p, \dots, \bar{x}_{P-1}^p, \bar{x}_0^p, \dots, \bar{x}_{P-1}^p, \dots).$$

The points on  $X_\phi^P(\Pi)$  will be denoted by  $x_k^\phi$ , i.e.  $x_k^\phi = \bar{x}_{(k+\phi) \bmod P}^p$ , and the corresponding control values by  $u_k^\phi$ .  $\square$

For any  $P$ -periodic trajectory, the ordered tuple  $(\bar{x}_0^p, \dots, \bar{x}_{P-1}^p) = (x_0^p, \dots, x_{P-1}^p)$  is a  $P$ -periodic orbit  $\Pi$ . Conversely, for every  $P$ -periodic orbit  $\Pi$  and any  $\phi \in \{0, \dots, P-1\}$  the sequence  $X_\phi^P(\Pi)$  from (ii) is a  $P$ -periodic trajectory in the sense of (i).

We extend the definition of (strict) dissipativity to periodic orbits as a generalization of [1]. To this end, in what follows we denote the particular periodic orbit for which the system is dissipative by  $\Pi^*$  with corresponding control sequence  $u^*$ . The corresponding elements will be denoted by  $\bar{x}_k^{p*}$  and  $\bar{u}_k^{p*}$ . Given a phase  $\phi$ , we denote the elements of the corresponding  $P$ -periodic trajectory  $X_\phi^P(\Pi^*)$  by  $(x_0^{\phi*}, x_1^{\phi*}, \dots)$  and the corresponding control values by  $(u_0^{\phi*}, u_1^{\phi*}, \dots)$ . Let us define for a point  $x$  and the periodic orbit  $\Pi^*$  the distance

$$|x|_{\Pi^*} := \min_{\bar{x}_k^{p*} \in \Pi} \|x - \bar{x}_k^{p*}\|, \quad (2)$$

and the distance

$$|(x, u)|_{\Pi_U^*} := \min_{(\bar{x}_k^{p*}, \bar{u}_k^{p*}) \in \Pi_U} \|x - \bar{x}_k^{p*}\| + \|u - \bar{u}_k^{p*}\|.$$

Let us define functions  $\sigma^\bullet(x, u)$ ,  $k \in \mathbb{N}_0$ , as

$$\sigma^A(x, u) := \rho(|(x, u)|_{\Pi_U^*}) \quad (3)$$

$$\text{or} \quad \sigma^B(x, u) := \rho(|x|_{\Pi^*}), \quad (4)$$

with  $\rho$  being a positive definite function. We remark that in case of (4) function  $\sigma^B(\cdot, \cdot)$  does not depend on  $u$ , but in order to obtain a uniform notation in what follows we always write  $\sigma^\bullet(x, u)$ .

*Definition 3.3 ( $P$ -Periodic (Strict) Dissipativity):* The system (1) is  $P$ -periodic dissipative on a set  $\mathbb{Z} \subset \mathbb{X} \times \mathbb{U}$  with respect to the supply rate  $\ell(x, u) - \ell(x_k^\phi, u_k^\phi)$  if there exists a feasible  $P$ -periodic orbit  $\Pi^*$  a phase  $\phi$  and bounded storage functions  $\lambda_0, \dots, \lambda_{P-1}, \lambda_P, \dots : X \rightarrow \mathbb{R}$ , with  $\lambda_{k+P} = \lambda_k$  such that the inequalities

$$L_k(x, u) := \ell(x, u) - \ell(x_k^\phi, u_k^\phi) + \lambda_k(x) - \lambda_{k+1}(f(x, u)) \geq 0 \quad (5)$$

hold for all  $(x, u) \in \mathbb{Z}$ , where  $x_k^\phi$  are the elements of the sequence  $X_\phi^P(\Pi)$  and all  $k = 0, 1, \dots$ . If, in addition, there exist functions of the form (3) or (4) such that

$$L_k(x, u) \geq \sigma^\bullet(x, u), \quad \bullet \in \{A, B\} \quad (6)$$

holds, then the system (1) satisfies  $P$ -periodic strictly dissipativity of type A or B, respectively, on  $\mathbb{Z}$ .  $\square$

It is easily seen that for (4) this definition is equivalent to Definition 2.1 in case  $P = 1$ . Moreover, for  $P > 1$ , (strict) dissipativity might hold for more than one phase  $\phi$ . While this can be restrictive if one is interested in the actual computation of  $L_k(x, u)$ , this does not constitute any problem for the theoretical results that we aim at establishing next, i.e. optimal  $P$ -periodic operation (uniform suboptimal non  $P$ -periodic operation), and sufficiency of strict  $P$ -periodic dissipativity for  $P$ -periodic stability of EMPC.

*Remark 3.4:* As it holds that  $|(x, u)|_{\Pi_U^*} \geq |x|_{\Pi^*}$ , Definition (6) in the sense A implies Definition (6) in the sense B.  $\square$

*Remark 3.5:* Note that, the time-varying and phase-dependent definition  $\sigma_k^C(x, u) := \rho(\|x - x_k^\phi\|)$  would at first look like the natural extension of the steady state case. However, in contrast to the time varying case in [13], this definition does not work in the time invariant setting of this paper. More precisely, if  $L_k(x, u) \geq \sigma_k^C(x, u)$  for phase  $\phi_1$  and the rotated cost of the  $P$ -periodic optimal trajectory

is evaluated for phase  $\phi_2 \neq \phi_1$ , then we obtain the inequality  $\sum_{k=0}^P L_k(x_k^{\phi_2}, u_k^{\phi_2}) \geq \rho(\|x_k^{\phi_2} - x_k^{\phi_1}\|)$ , which can never be satisfied since  $\sum_{k=0}^P L_k(x_k^{\phi_2}, u_k^{\phi_2}) = 0$  and  $\rho(\|x_k^{\phi_2} - x_k^{\phi_1}\|) > 0$ .  $\square$

#### A. Optimal $P$ -periodic Operation and (Strict) Dissipativity

A  $P$ -periodic orbit  $\Pi^*$  with corresponding control sequence  $u^*$  is called *optimal* if it has minimal period  $P^*$  and corresponds to the state-control pairs  $\Pi_U^*$  defined as

$$(P^*, \Pi_U^*) \in \operatorname{argmin}_{P, \Pi_U} \frac{1}{P} \sum_{k=0}^{P-1} \ell(x_k, u_k), \quad (7)$$

where minimization is carried out over all periods  $P \geq 1$  and all periodic state-control sequences  $\Pi_U$  of minimal period  $P$ . We emphasise that, in general, the argmin is not unique. Also note that the minimum might not exist.

The average optimal  $P$ -periodic cost (which is independent of  $\phi$ ) is given by

$$\ell_P^* := \frac{1}{P} \sum_{k=0}^{P-1} \ell(x_k^{\phi^*}, u_k^{\phi^*}).$$

For a real vector valued sequence  $v = (v_0, v_1, \dots)$  we define the set of  $P$ -step asymptotic averages as

$$\operatorname{Av}^P[v] = \{\bar{v} \in \mathbb{R}^{n_v} : \exists t_n \rightarrow +\infty : \lim_{n \rightarrow \infty} \frac{\sum_{k=0}^{t_n} \sum_{j=0}^{P-1} v_{Pk+j}}{P(t_n+1)} = \bar{v}\},$$

noting that this set is actually independent of  $P$  if the sequence  $v_k$  is bounded.

Let us now define, analogously to [2] and [9], several optimal  $P$ -periodic operation concepts.

**Definition 3.6 (Optimal  $P$ -Periodic Operation):** The system (1) is *optimally  $P$ -periodically operated* at a periodic orbit  $\Pi^*$  with respect to the stage cost  $\ell$ , if for each solution satisfying  $(x_k, u_k) \in \mathbb{Z}$  for all  $k = 0, 1, \dots$ , the following holds:

$$\operatorname{Av}^P[\ell(x, u)] \subset [\ell_P^*, \infty). \quad (8)$$

**Definition 3.7 (Suboptimal non  $P$ -Periodic Operation):** The system (1) is *suboptimally non  $P$ -periodically operated* at a periodic orbit  $\Pi^*$  with respect to the stage cost  $\ell$  and the functions  $\sigma^\bullet$  from (3) or (4), if it is optimally  $P$ -periodically operated and in addition one of the following two conditions holds:

$$\operatorname{Av}^P[\ell(x, u)] \subset (\ell_P^*, \infty), \quad (9a)$$

$$\text{there is } \phi \in \{0, \dots, P-1\} \text{ with } \liminf_{k \rightarrow \infty} \sigma^\bullet(x_k, u_k) = 0. \quad (9b)$$

**Definition 3.8 (Uniform Suboptimal non  $P$ -Periodic Operation):** The system (1) is *uniformly suboptimally non  $P$ -periodically operated* at a periodic orbit  $\Pi^*$  with respect to the stage cost  $\ell$  and the functions  $\sigma^\bullet$  from (3) or (4), if it is suboptimally non  $P$ -periodically operated and in addition for each  $\delta > 0$  there exists an integer  $\bar{t} \geq 1$  such that one of the following two conditions holds:

$$\sum_{k=0}^t \sum_{j=0}^{P-1} \frac{\ell(x_{Pk+j}, u_{Pk+j})}{Pt} \geq \ell_P^*, \text{ for all } t \geq \bar{t}, \quad (10a)$$

$$\text{there is } \phi \in \{0, \dots, P-1\} \text{ with } \sigma^\bullet(x_k, u_k) \leq \delta, \text{ for } P \text{ consecutive } k \in [1, \bar{t}]. \quad (10b)$$

**Remark 3.9:** We note that the actual behavior of the trajectories satisfying (10b) differs depending on  $\sigma^\bullet$ .

In case of (3), if Property (10b) holds for sufficiently small  $\delta$ , then by continuity of  $f$  from  $\rho(\|x_k^{\phi^*} - \bar{x}_k^P\| + \|u_k^{\phi^*} - \bar{u}_k^P\|) \leq \delta$  we obtain  $f(x_k^{\phi^*}, u_k^{\phi^*}) \approx \bar{x}_{k^+}^{P^*}$  with  $k^+ = k + 1 \pmod{P^*}$ . Since the periodic orbit consists of finitely many distinct points, for sufficiently small  $\delta > 0$  this implies  $\sigma^A(f(x_k^{\phi^*}, u_k^{\phi^*}), u_{k^+}^{\phi^*}) \geq \rho(\|f(x_k^{\phi^*}, u_k^{\phi^*}) - \bar{x}_{k^+}^{P^*}\|) > \delta$  for all  $j \neq k^+$  which together with  $\sigma^A(f(x_k^{\phi^*}, u_k^{\phi^*}), u_{k^+}^{\phi^*}) \leq \delta$  yields  $\rho(\|f(x_k^{\phi^*}, u_k^{\phi^*}) - \bar{x}_{k^+}^{P^*}\| + \|u_{k^+}^{\phi^*} - \bar{u}_{k^+}^{P^*}\|) \leq \delta$ . As a consequence, any state-control sequence sufficiently close to  $\Pi_U$  approximately follows the periodic motion.

In contrast to this, in case of (4) we can only conclude that the solution stays near the set  $\Pi_U$  but it need not approximately follow the periodic motion. While it is possible to re-establish approximate periodicity in case  $\Pi_U^*$  is the unique minimiser of  $J_P(x, u)$  over all (not necessarily periodic) orbits of length  $P$ , this will require additional arguments in the subsequent proofs and does not directly follow from (4), see also Remark 4.7.  $\square$

We can now state the following theorems relating dissipativity and optimal operation of the system.

**Theorem 3.10:** Assume that system (1) is (strictly)  $P$ -periodically dissipative on  $\mathbb{Z}$  with respect to the supply rate  $\ell(x_k, u_k) - \ell(x_k^*, u_k^*)$  and  $\sigma^\bullet$  from (3) or (4). Then system (1) is optimally  $P$ -periodically operated (uniformly suboptimally non  $P$ -periodically operated) at the optimal  $P$ -periodic trajectory  $X_\phi^P(\Pi^*)$ .  $\square$

*Proof:* The proof follows with appropriate adaptations from the one given in [2, Proposition 6.4] and [9, Theorem 1] for the case  $P = 1$ . We have

$$\begin{aligned} 0 &= \lim_{T \rightarrow \infty} \frac{\lambda_{PT}(x_{PT}) - \lambda_0(x_0)}{PT} \\ &= \lim_{T \rightarrow \infty} \sum_{k=0}^{T-1} \sum_{j=0}^{P-1} \frac{\lambda_{Pk+j+1}(x_{Pk+j+1}) - \lambda_{Pk+j}(x_{Pk+j})}{PT} \\ &\leq \liminf_{T \rightarrow \infty} \sum_{k=0}^{T-1} \sum_{j=0}^{P-1} \frac{\ell(x_{Pk+j}, u_{Pk+j})}{PT} - \ell_P^*. \end{aligned}$$

This establishes the first claim. If strict  $P$ -periodic dissipativity holds, then there is a phase  $\phi$  with

$$\begin{aligned} 0 &\leq \liminf_{T \rightarrow \infty} \sum_{k=0}^{T-1} \sum_{j=0}^{P-1} \frac{\sigma^\bullet(x_{Pk+j}, u_{Pk+j})}{PT} \\ &\leq \liminf_{T \rightarrow \infty} \sum_{k=0}^{T-1} \sum_{j=0}^{P-1} \frac{\ell(x_{Pk+j}, u_{Pk+j})}{PT} - \ell_P^*, \end{aligned}$$

and two cases are possible:

- 1)  $\liminf_{T \rightarrow \infty} \sum_{k=0}^{T-1} \sum_{j=0}^{P-1} \frac{\ell(x_{Pk+j}, u_{Pk+j})}{PT} > \ell_P^*$ , which implies  $\operatorname{Av}^P[\ell(x, u)] \subset (\ell_P^*, \infty)$ , or
- 2)  $\liminf_{T \rightarrow \infty} \sum_{k=0}^{T-1} \sum_{j=0}^{P-1} \frac{\ell(x_{Pk+j}, u_{Pk+j})}{PT} = \ell_P^*$ , hence  $\liminf_{k \rightarrow \infty} \sigma^\bullet(x_k, u_k) = 0$ .

This proves that strict  $P$ -periodic dissipativity entails suboptimal non  $P$ -periodic operation. It remains to prove uniform suboptimal non  $P$ -periodic operation.

For each feasible solution and  $t \geq 0$ , (6) and boundedness of the  $\lambda_k$  entails that

$$-c := -2 \sup_{\substack{0 \leq k \leq P-1 \\ x \in \mathbb{X}_Z}} |\lambda_k(x)| \leq \lambda_{Pt}(x_{Pt}) - \lambda_0(x_0)$$

$$\leq \sum_{k=0}^{t-1} \sum_{j=0}^{P-1} [\ell(x_{Pk+j}, u_{Pk+j}) - \sigma^\bullet(x_{Pk+j}, u_{Pk+j})] - Pt\ell_P^*.$$

Let  $\delta > 0$  be fixed and choose  $\bar{t} := \left\lceil \frac{c}{\rho(\delta)} \right\rceil + 1$ . Then two cases are possible:

- 1)  $\sum_{k=0}^{t-1} \sum_{j=0}^{P-1} \ell(x_{Pk+j}, u_{Pk+j}) > Pt\ell_P^*$  for all  $t \geq \bar{t}$ , or
- 2)  $\sum_{j=0}^{P-1} \sigma(x_{Pk+j}, u_{Pk+j}) \leq c/\bar{t}$  for at least one  $k \in [1, \bar{t}]$ ,  
implying  $\sigma^\bullet(x_j, u_j) \leq c/\bar{t}$  and thus  $\sigma^\bullet(x_j, u_j) \leq \delta$  for  $j = Pk, \dots, (P+1)k-1$

which concludes the proof.  $\square$

#### IV. PERIODIC STABILITY OF ECONOMIC MPC

Let us consider the following MPC problem

$$V^i(x) = \min_{x_0, u_0, \dots, x_N} J^i(x, u) \quad (11a)$$

$$\text{with } J^i(x, u) = \sum_{k=0}^{N-1} \ell(x_k, u_k) + V_f^{N+i}(x_N) \quad (11b)$$

$$\text{s.t. } x_0 = x, \quad (11c)$$

$$x_{k+1} = f(x_k, u_k), \quad (11d)$$

$$(x_k, u_k) \in \mathbb{Z}, \quad (11e)$$

$$x_N \in \mathbb{X}_f^{N+i}, \quad (11f)$$

where the (periodic) terminal set and cost depend on the current time instant  $i$ . We note that this time-dependence can be used in order to induce a fixed phase for the EMPC closed loop trajectory. Note however that even if a terminal constraint with phase  $\phi_1$  is introduced, the closed-loop solution can in general have a phase  $\phi_2 \neq \phi_1$ . We may also use terminal costs and constraints which are independent of  $i$ , in which case the phase is not fixed. We also remark that for non constant  $\mathbb{X}_f^{N+i}$  the feasible sets  $\mathbb{X}_N^i$ , i.e., the sets of all  $x$  for which the constraints in (11) can be satisfied, depend periodically on  $i$ .

Let us introduce the following assumptions.

*Assumption 4.1:* The sets  $\mathbb{X}_f^{N+i}$  are compact.  $\square$

*Assumption 4.2:* The stage cost  $\ell(\cdot, \cdot)$  and system dynamics  $f(\cdot, \cdot)$  are continuous on  $\mathbb{Z}$ . The terminal cost function  $V_f^{N+i}(\cdot)$  is continuous on the terminal region  $\mathbb{X}_f^{N+i}$ .  $\square$

*Assumption 4.3 (P-Periodic Strict Dissipativity):* System (1) is strictly dissipative at a periodic orbit  $\Pi^*$  with respect to the supply rate  $\ell(x, u) - \ell(x_k^{\phi^*}, u_k^{\phi^*})$  and  $\sigma^\bullet$  from (3) or (4). Moreover, the storage function  $\lambda(\cdot)$  is bounded and continuous in every point  $x^{P^*} \in \Pi^*$ .  $\square$

*Assumption 4.4:* The value function  $V^i(\cdot)$  is bounded on  $\mathbb{X}_N^i$  and continuous in every point  $x^{P^*} \in \Pi^*$ .  $\square$

Let us define the rotated MPC problem and the corresponding rotated value function as

$$\bar{V}^i(x) = \min_{x_0, u_0, \dots, x_N} \bar{J}^i(x, u) \quad (12a)$$

$$\text{with } \bar{J}^i(x, u) = \sum_{k=0}^{N-1} L_{k+i}(x_k, u_k) + \bar{V}_f^{N+i}(x_N) \quad (12b)$$

$$\text{s.t. } x_0 = x, \quad (12c)$$

$$x_{k+1} = f(x_k, u_k), \quad (12d)$$

$$(x_k, u_k) \in \mathbb{Z}, \quad (12e)$$

$$x_N \in \mathbb{X}_f^{N+i}, \quad (12f)$$

where the rotated terminal and stage cost are phase-dependent and defined respectively as  $\bar{V}_f^{N+i}(x) := V_f^{N+i}(x) + \lambda_{N+i}(x)$  and  $L_k = L_{k(\text{mod } P^*)}$  from (5). These definitions imply

$$\bar{J}_N^i(x, u) = J_N^i(x, u) + \lambda_i(x) - \sum_{k=0}^{N-1} \ell(x_{k+i}^{\phi^*}, u_{k+i}^{\phi^*}) \quad (13)$$

and thus the rotated MPC problem delivers the same optimal trajectories and control sequences as the original problem, see also [13].

Let us consider a family of periodic terminal regions  $\mathbb{X}_f^k \subset \mathbb{X}$  and terminal costs  $V_f^k$  satisfying the following assumptions.

*Assumption 4.5:* The terminal regions are periodic, i.e.,  $\mathbb{X}_f^{k+P} = \mathbb{X}_f^k$  for all  $k \geq 0$  and the periodic terminal regions  $\mathbb{X}_f^k$  contain the states  $x_k^{\phi^*}$  of the periodic trajectory  $X_\phi^P(\Pi^*)$  from Assumption 4.3. Moreover, the terminal costs are periodic, i.e.,  $V_f^{k+P} = V_f^k$  for all  $k \geq 0$  and there exist a terminal control law  $\kappa_f^k : \mathbb{X}_f^k \rightarrow \mathbb{U}$  with  $\kappa_f^{k+P} = \kappa_f^k$  for all  $k \geq 0$  such that, at a given time instant  $i$ , for all  $x \in \mathbb{X}_f^{N+i}$  the inclusion  $f(x, \kappa_f^{N+i}(x)) \in \mathbb{X}_f^{N+i+1}$  and the inequality

$$V_f^{N+i+1}(f(x, \kappa_f^{N+i}(x))) \leq V_f^{N+i}(x) - \ell(x, \kappa_f^{N+i}(x)) + \ell(x_{N+i}^{\phi^*}, u_{N+i}^{\phi^*}),$$

holds.  $\square$

We remark that in case  $\mathbb{X}_f^k = \{x_k^{\phi^*}\}$ , Assumption 4.5 is satisfied with  $\kappa_f^k \equiv u_k^{\phi^*}$  and  $V_f^k \equiv 0$ . The simplest example for time invariant terminal conditions are  $\mathbb{X}_f = \{x^{P^*} \in \Pi^*\}$  with  $\kappa_f(x_k^{P^*}) = u_k^{P^*}$  and again  $V_f \equiv 0$ . We also note that Assumption 4.5 is satisfied for the original MPC problem if and only if it is satisfied for the rotated problem. For an analysis of a periodic EMPC scheme without any terminal conditions we refer to [10].

*Theorem 4.6:* Let Assumptions 4.1, 4.2, 4.3, 4.4 and 4.5 hold. Then the rotated optimal value function  $\bar{V}(x)$  is a Lyapunov function and, for  $\sigma^A$ , i.e. from (3), there exists a phase  $\phi$  such that the trajectory  $x_k^{\phi^*}$  corresponding to the optimal periodic orbit  $\Pi^*$  is asymptotically stable for the closed loop system. For  $\sigma^B$ , i.e. from (4), the optimal periodic orbit  $\Pi^*$  is an asymptotically stable set for the closed loop system.

*Proof:* The proof uses ideas similar to the steady state case [1] with appropriate adaptations. We define  $\sigma^*(x) := \inf_{u \in U} \sigma^A(x, u) = \inf_{u \in U} \sigma^B(x, u)$ . Formula (13) and the boundedness and continuity in every  $x_k^{P^*}$  from  $\Pi^*$  of  $V^i$  and  $\lambda$  ensured by Assumptions (4.3) and (4.4) imply that  $\bar{V}^i$  is also bounded on  $\mathbb{X}_N^i$  and continuous in every  $x_k^{P^*}$  from  $\Pi^*$ . Moreover, strict dissipativity implies  $\bar{V}^i(x) \geq L_k(x, u) \geq \sigma^*(x)$  and  $L_k(x_k^{P^*}, u_k^{P^*}) = 0$  implies  $\bar{V}^i(x_k^{P^*}) = 0$ . Together, these properties ensure the existence of  $\mathcal{K}$  functions  $\hat{\alpha}$  and  $\alpha$  such that

$$\hat{\alpha}(\sigma^*(x)) \leq \bar{V}^i(x) \leq \alpha(\sigma^*(x)).$$

Note that local loss of controllability near the periodic optimal trajectory can entail a discontinuity of  $V^i(\cdot)$  at  $\Pi^*$  and hence of  $\bar{V}^i(\cdot)$ . If the cost  $\bar{V}^i(\cdot)$  is not continuous at the periodic optimal trajectory, we cannot establish the upper bound  $\bar{V}^i(x) \leq \alpha(\sigma^*(x))$ .

In order to prove descent of the rotated value function  $\bar{V}^i(x)$ , let us define the optimal state and control trajectory as

$$X_i^{\text{MPC}} = (x_{0,i}^{\text{MPC}}, x_{1,i}^{\text{MPC}}, \dots, x_{N,i}^{\text{MPC}}),$$

$$U_i^{\text{MPC}} = (u_{0,i}^{\text{MPC}}, u_{1,i}^{\text{MPC}}, \dots, u_{N-1,i}^{\text{MPC}}).$$

Let us moreover define a feasible candidate trajectory for the MPC problem at the next time step as

$$\bar{X}_{i+1} = (x_{1,i}^{\text{MPC}}, \dots, x_{N,i}^{\text{MPC}}, f(x_{N,i}^{\text{MPC}}, \kappa_{f,i}^k(x_{N,i}^{\text{MPC}}))),$$

$$\bar{U}_{i+1} = (u_{1,i}^{\text{MPC}}, \dots, u_{N-1,i}^{\text{MPC}}, \kappa_{f,i}^k(x_{N,i}^{\text{MPC}})).$$

The rotated objective value associated with this trajectory is given by

$$\begin{aligned} \bar{J}^{i+1}(x_{1,i}^{\text{MPC}}) &= \bar{V}^i(x_{0,i}^{\text{MPC}}) - L_i(x_{0,i}^{\text{MPC}}, u_{0,i}^{\text{MPC}}) \\ &\quad - V_f^{N+i}(x_{N,i}^{\text{MPC}}) + V_f^{N+i}(f(x_{N,i}^{\text{MPC}}, \kappa_f^{N+i}(x_{N,i}^{\text{MPC}}))) \\ &\quad + \ell(x_N, \kappa_f^{N+i}(x_N)) - \ell(x_{N+1+i}^{\phi^*}, u_{N+1+i}^{\phi^*}) \\ &\leq \bar{V}^i(x_{0,i}^{\text{MPC}}) - L_i(x_{0,i}^{\text{MPC}}, u_{0,i}^{\text{MPC}}). \end{aligned}$$

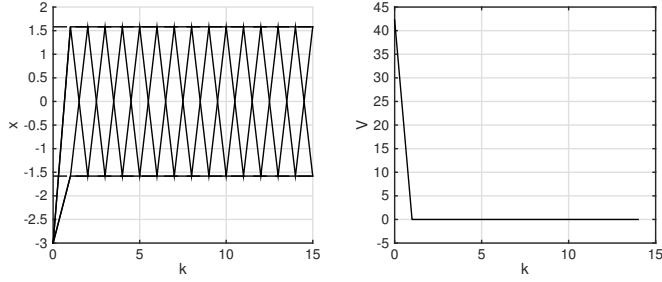


Fig. 1. Example 5.1. Left graph: closed-loop trajectory (continuous line) obtained starting from  $\hat{x}_0 = -3$ . The periodic optimal states are displayed in dotted line. Right graph: Value function of the rotated MPC problem.

Optimality implies  $\bar{V}^{i+1}(x_{1,i}^{\text{MPC}}) \leq \bar{J}^{i+1}(x_{1,i}^{\text{MPC}})$  and hence

$$\begin{aligned} \bar{V}^{i+1}(x_{1,i}^{\text{MPC}}) - \bar{V}^i(x_{0,i}^{\text{MPC}}) &\leq -L_i(x_{0,i}^{\text{MPC}}, u_{0,i}^{\text{MPC}}) \\ &\leq -\sigma^\bullet(x_{0,i}^{\text{MPC}}, u_{0,i}^{\text{MPC}}). \end{aligned}$$

The periodic family of rotated value functions is hence a family of Lyapunov functions for the nonlinear system; particularly,  $\bar{V}^i$  converges to 0 along the closed loop trajectory. From this, for  $\sigma^B$  from (4), the claimed stability properties immediately follow.

For  $\sigma^A$  from (3), the lower bound  $\hat{\alpha}(\sigma^*(x))$  of the Lyapunov functions only implies the convergence of the states of the closed loop to  $\Pi^*$  but not necessarily of the controls. Hence, the proof so far only shows asymptotic stability of the set  $\Pi^*$  but not of the periodic trajectory  $x_k^{\phi^*}$  corresponding to  $\Pi^*$ . However, from the last inequality, above, we obtain

$$\sigma^A(x_{0,i}^{\text{MPC}}, u_{0,i}^{\text{MPC}}) \leq \bar{V}^i(x_{0,i}^{\text{MPC}})$$

implying that since  $V^i$  tends to 0 the value  $\sigma^A(x_{0,i}^{\text{MPC}}, u_{0,i}^{\text{MPC}})$  also tends to 0. By (3) this yields that  $|(x_{0,i}^{\text{MPC}}, u_{0,i}^{\text{MPC}})|_{\Pi_u}$  tends to 0 as  $V_i$  tends to 0 and thus asymptotic stability of the trajectory corresponding to the periodic orbit  $\Pi^*$  follows by similar arguments as in Remark 3.9.  $\square$

*Remark 4.7:* In case of strict dissipativity of type B, i.e. with  $\sigma^B$  from (4), asymptotic stability of the periodic trajectory  $x_k^{\phi^*}$  follows if the optimal periodic orbit  $\Pi^*$  is the unique minimiser of  $J_P(x, u)$  over all (not necessarily periodic) orbits of length  $P$ . Indeed, in this case for  $x_{0,i}^{\text{MPC}}$  sufficiently close to  $\Pi^*$ , due to continuity  $X_i^{\text{MPC}}$  must approximately follow  $x_k^{\phi^*}$  because otherwise we would obtain a contradiction to the optimality of  $X_i^{\text{MPC}}$ .  $\square$

## V. EXAMPLES

The following examples illustrate the proposed concepts.

*Example 5.1 (Strict Dissipativity of type B):*

Consider the 1d nonlinear dynamics  $f(x, u) = -x + u$  and stage cost

$$\ell(x, u) = (x - 2)(x - 1)(x + 1)(x + 2).$$

The optimal trajectory can either be of period  $P = 1$ , i.e. one of the two steady states  $x_s^{1,2} = \pm \frac{\sqrt{10}}{2}$ , or of period  $P = 2$ , with  $\Pi^* = \left(\frac{\sqrt{10}}{2}, -\frac{\sqrt{10}}{2}\right)$  and  $u_1^{p*} = u_0^{p*} = 0$ . Using  $\lambda_0(x) = \lambda_1(x) = 0$ , it can be verified that  $L_0(x, u) = L_1(x, u)$  satisfy the strict dissipation of type B, i.e. with  $\sigma^B(\cdot, \cdot)$  from (4).

As the control is not bounded and it does not enter the cost, any MPC scheme will stabilise the system in one step. The solution is however not unique and we can conclude from Definition (4) that the system will be stabilised to the set of states which includes the periodic optimal trajectory. However, both staying at one of the

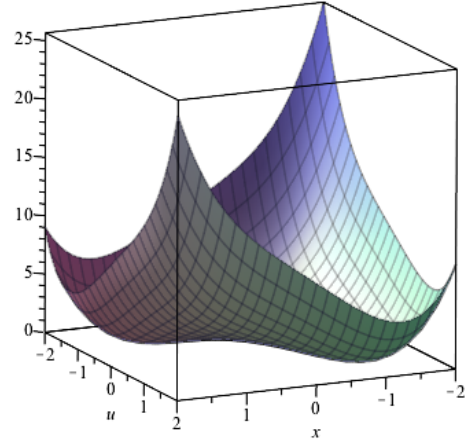


Fig. 2. Example 5.2: graph of the rotated stage costs  $L_1 = L_2$ .

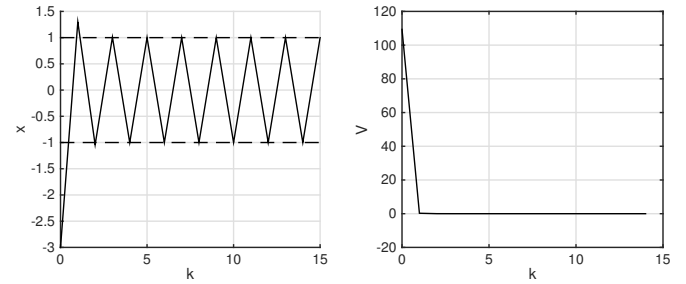


Fig. 3. Example 5.2. Left graph: closed-loop trajectory (continuous line) obtained starting from  $\hat{x}_0 = -3$ . The periodic optimal states are displayed in dotted line. Right graph: Value function of the rotated MPC problem.

steady states and moving from one to the other are optimal moves. Using the initial condition  $\hat{x}_0 = -3$  and the terminal constraint  $x_N = (-\sqrt{10}/2)^{i+N+1}$ , all possible closed-loop trajectories and the value of the rotated problem are displayed in Figure 1. Note that, as the control effort is not penalised by the stage cost, the system is stabilised to the optimal operation in one step.  $\square$

*Example 5.2 (Strict Dissipativity of type A):* Consider the 1d system with dynamics  $f(x, u) = u$  and stage cost

$$\ell(x, u) = x^4 - \frac{x^3}{3} - 2x^2 + x + (x + u)^2.$$

Using

$$\lambda_1(x) = -\frac{x^4}{2} + \frac{x^3}{6} + x^2 - \frac{x}{2} + \frac{2}{3}, \quad \lambda_2(x) = -\frac{x^4}{2} + \frac{x^3}{6} + x^2 - \frac{x}{2},$$

$\Pi = (1, -1)$  and  $\Pi_U = ((1, -1), (-1, 1))$ , one obtains

$$L_1(x, u) = L_2(x, u) = \frac{x^4}{2} - \frac{x^3}{6} + \frac{x}{2} + 1 + \frac{u^4}{2} - \frac{u^3}{6} + \frac{u}{2} + 2xu.$$

One checks that this polynomial has exactly two local minima at  $(1, -1)$  and  $(-1, 1)$  at which its value is 0, cf. Figure 2. Hence, it is positive elsewhere and since it grows unboundedly for  $|x|, |u| \rightarrow \infty$ , we can find  $\rho \in \mathcal{K}_\infty$  such that (6) holds with  $\sigma^\bullet(\cdot, \cdot)$ ,  $\bullet = \{A, B\}$ , i.e. from (3) or (4).

For an MPC scheme with horizon  $N = 5$ , initial condition  $\hat{x}_0 = -3$  and terminal constraint  $x_N = (-1)^{i+N+1}$ , the closed-loop trajectory and the value of the rotated problem are displayed in Figure 3.  $\square$

*Example 5.3 (Strict Dissipativity of type B for a 2d system):* Define

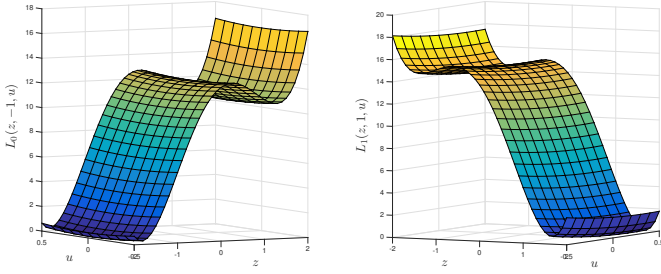


Fig. 4. Example 5.3: Rotated stage cost  $L_0(z, y, u)$  with  $y = -1$  fixed (left) and  $L_1(z, y, u)$  with  $y = 1$  fixed (right).

$x = [z, y]^T$  and consider the nonlinear dynamics and stage cost

$$f(x, u) = \begin{bmatrix} -0.9z + yu \\ -y \end{bmatrix},$$

$$\ell(x, u) = (z - 1.9)(z - 0.9)(z + 1.1)(z + 2.1) + (u - 20)^2,$$

with constraint  $y \in \{-1, 1\}$ .

Let us consider the case  $\{y_0^{p*}, y_1^{p*}\} = \{1, -1\}$ . The optimal trajectory is periodic with period  $P = 2$  and can be computed numerically:  $\Pi^* = \{(z_0^{p*}, y_0^{p*})^T, (z_1^{p*}, y_1^{p*})^T\}$  with  $z_0^{p*} \approx -1.8294$ ,  $z_1^{p*} \approx 1.6719$ ,  $\{u_0^{p*}, u_1^{p*}\} \approx \{0.0254, 0.3247\}$  and Lagrange multipliers associated to the  $z$ -variable of the dynamic constraints  $\{\lambda_0^p, \lambda_1^p\} \approx \{39.9492, -39.3506\}$ .

Using

$$\lambda_0(x) = (dy - c)(z - z_0) + (1 - y)e \quad \lambda_1(x) = (dy - c)(z - z_1),$$

with  $d = \frac{\lambda_0^p - \lambda_1^p}{2}$  and  $c = \frac{\lambda_0^p + \lambda_1^p}{2}$ , we obtain

$$L_0(x, u) = \ell(x, u) - \ell(x_0, u_0) + \lambda_0(x) - \lambda_1(f(x, u)) + (1 - y)e,$$

$$L_1(x, u) = \ell(x, u) - \ell(x_1, u_1) + \lambda_1(x) - \lambda_0(f(x, u)) - (1 + y)e.$$

Computing the minima of these functions reveals that for  $e \approx 6.89763344$  the functions  $L_k$  satisfy the strict dissipation inequalities (6) for  $\sigma^A(\cdot, \cdot)$ , cf. Figure 4.

For an MPC scheme with horizon  $N = 5$  and an initial condition  $\hat{x}_0 = (-3, 1)$  with terminal constraint  $x_N = (-1)^{i+N+1} (z_0^p, -1)$ , the closed-loop trajectory and the value of the rotated problem are displayed in Figure 5.  $\square$

## VI. DISCUSSION AND CONCLUSIONS

In this paper, we have presented an extension of strict dissipativity to periodic systems. We have proven that several previous results obtained for the steady state case extend to our setting for periodic operation. These theoretical results have been illustrated using several numerical examples.

The proposed setting straightforwardly extends to the case of multistep MPC [6, Section 7.4]. Future work will include a multistep version of our results, as well as the extending them to the case of MPC schemes without a terminal constraint nor cost.

The major limitations of the current stability theory, both for the steady state and the periodic case, include the following:

- 1) while sufficiency of strict dissipativity for stability has been proven in [1] and in the current paper, to the authors' knowledge, no result on its necessity has been obtained yet
- 2) in general it can be very difficult to prove the existence of a storage function which satisfies the strict dissipativity condition
- 3) the storage function is assumed to be bounded and continuous in  $x_k^{p*}$  from  $\Pi^*$ .
- 4) In each of our examples, the functions  $L_k$  are identical for all  $k$ . So far we were not able to determine whether this is just

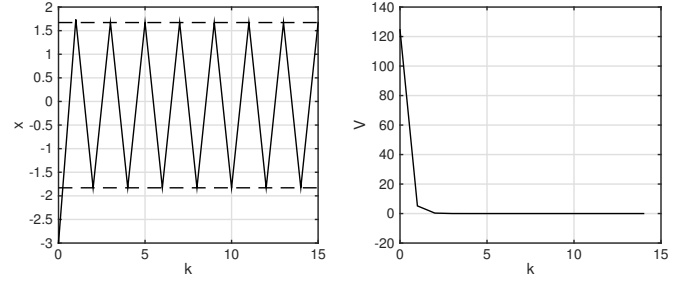


Fig. 5. Example 5.3. Left graph: closed-loop trajectory (continuous line) obtained starting from  $\hat{x}_0 = -3$ . The periodic optimal states are displayed in dotted line. Right graph: Value function of the rotated MPC problem.

a coincidence or whether there is a systematic reason for this fact.

Future research will aim at developing the theory further so as to overcome these limitations.

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